Simulating qubit correlations with classical communication

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Qubits are better!

$0\mapsto |0 angle, 1\mapsto |1 angle$

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Qubits are strictly better!!

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$$x_0, x_1 \in \{0, 1\}$$

$$y \in \{0, 1\}$$

$$m \in \{0, 1\}$$

$$b \in \{0, 1\}$$

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win if
$$b = xy$$

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win if
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 $p_{\text{classic}} \le \frac{3}{4}$

$$x_0, x_1 \in \{0, 1\}$$

$$y \in \{0, 1\}$$

$$|\psi\rangle \in \mathcal{H}_2$$

$$b \in \{0, 1\}$$

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$$\begin{array}{c} x_0, x_1 \in \{0,1\} \\ & y \in \{0,1\} \\ & |\psi\rangle \in \mathcal{H}_2 \\ & b \in \{0,1\} \\ & \text{win if } b = xy \end{array}$$

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$$p_{\text{quantum}} \le \frac{2+\sqrt{2}}{4} \approx 85\%$$

What if Alice sends 2 bits?

Prepare-and-Measure







 $b \in \{0, \dots, N\}$



 $|\psi\rangle\in\mathcal{H}_2$



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Quantum teleportation





 $|\psi\rangle\in\mathcal{H}_2$



 $b \in \{0, \dots, N\}$

No extra resource?







 $b \in \{0, \dots, N\}$



 $|\psi\rangle\in\mathcal{H}_2$



Qubit simulation requires unlimited shared randomness



 $m \in \{0, \ldots, d_C - 1\}$



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Massar, Bacon, Cerf, and Cleve, PRA (2001)

Prepare and Measure with Shared Randomness



The problem:





 $|\psi\rangle \in \mathcal{H}_2$



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- How about 1trit+SR?
- Buhrman, Cleve, Massar, de Wolf, Rev. Mod. Phys. (2010). Non-locality and communication complexity Many results, but not much about minimal worst case scenarios...

Our goal:

▶ 1: Analyse the trit vs Qubit case in detail

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- ▶ 1: Analyse the trit vs Qubit case in detail
- 2: Understand the power and limitations of POVMs

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For some tasks, a trit is better than qubit

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 $x_0, x_1 \in \{0, 1\}$



$$m \in \{0, 1, 2\}$$

 $y \in \{0,1\}$



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win if
$$b = xy$$

 $p_{\text{trit}} \le \frac{7}{8}$

For some tasks, a trit is better than qubit (Holevo bound!)

 $x \in \{0, 1, 2\}$



 $m \in \{0, 1, 2\}$



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 $|\psi\rangle\in\mathcal{H}_2$



1: Question?

Are trits strictly better than qubits?

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RESULT 1

Are trits strictly better than qubits? No!



For some tasks, a qubit is better than trit



RESULT 1

$$|\psi_m\rangle \in \checkmark$$

$$|\psi\rangle \in \mathcal{H}_2$$

$$\psi \in \{0,1\}$$

$$\operatorname{prob}\left(b \mid |\psi_m\rangle, M_y\right) = \operatorname{Tr}\left(|\psi_m\rangle\langle\psi_m| M_b|_y\right)$$

RESULT 1

$$x \in \{0, \dots, 5\}$$

$$m \in \{0, 1, 2\}$$

$$w \in \{0, 1, 2\}$$

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Recognise that the problem is a linear program (even with robustness considered)

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Optimisation trick to reduce complexity

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- Optimisation trick to reduce complexity
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- Extract a computer-assisted proof (as in Bavaresco, Murao, Quintino, PRL 127, 200504 (2021))
- ► Various examples, minimal: 6 preparations, 11 measurements

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Are 2bits strictly better than 1qubit?

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Are 2bits strictly better than 1qubit?





RESULT 2

2bits+SR is strictly better than qubits!



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Proof: Explicit recipe for classical simulation



use the Bloch sphere to states and POVM elements



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• $\lambda := (\vec{\lambda_1}, \vec{\lambda_2})$, random vectors on the sphere

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λ := (λ₁, λ₂), random vectors on the sphere
Instead of ρ = ½(I + x · σ), Alice sends c₁ = H(x · λ₁) and c₂ = H(x · λ₂) Heaviside: H(x) = 1 if x ≥ 0, H(x) = 0 if x < 0).

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• $\lambda := (\vec{\lambda_1}, \vec{\lambda_2})$, random vectors on the sphere

▶ Instead of $\rho = \frac{1}{2}(I + \vec{x} \cdot \vec{\sigma})$, Alice sends $c_1 = H(\vec{x} \cdot \vec{\lambda}_1)$ and $c_2 = H(\vec{x} \cdot \vec{\lambda}_2)$ Heaviside: H(x) = 1 if $x \ge 0$, H(x) = 0 if x < 0).

▶ Bob finds the Bloch vectors for the POVM elements, $B_b = p_b(I + \vec{y_b} \cdot \vec{\sigma})$ then sets $\vec{\lambda} := (-1)^{1+c_1} \vec{\lambda_1}$ when $|\vec{\lambda}_1 \cdot \vec{y}_b| \ge |\vec{\lambda}_2 \cdot \vec{y}_b|$ and $\vec{\lambda} := (-1)^{1+c_2} \vec{\lambda}_2$ otherwise.

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Finally, Bob outputs *b* with probability:

$$p(b|\{\vec{y}_b\}_b, \lambda) = \frac{p_b \Theta(\vec{y}_b \cdot \vec{\lambda})}{\sum_{j=1}^n p_j \Theta(\vec{y}_j \cdot \vec{\lambda})}$$
$$\Theta(x) := \begin{cases} x & \text{if } x \ge 0\\ 0 & \text{if } x < 0 \end{cases}.$$



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Why it works?

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Why it works? Well...

 $\forall \text{ qubit } \rho, \forall \text{ POVM } \{M_b\}$ $\int_{\lambda} d\lambda \pi(\lambda) \sum_{c=1}^{4} p_A(c|\rho, \lambda) p_B(b|\{M_b\}, c, \lambda) = \text{tr}(\rho M_b)$

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Why it works?

Lemma 1. Given two normalized vectors $\vec{x}, \vec{y} \in \mathbb{R}^3$ on the unit sphere S_2 , it holds that:

$$\frac{1}{\pi} \int_{S_2} H(\vec{x} \cdot \vec{\lambda}) \cdot \ \Theta(\vec{y} \cdot \vec{\lambda}) \ \mathrm{d}\vec{\lambda} = \frac{1}{2} (1 + \vec{x} \cdot \vec{y}) \,,$$

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where H(z) is the Heaviside function $(H(z) = 1 \text{ if } z \ge 0 \text{ and } H(z) = 0 \text{ if } z < 0)$ and $\Theta(z) := H(z) \cdot z$.

RESULT 2 extra

The fraction of rounds in which Alice is communicating only a single bit to Bob has measure zero.

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- This holds for any protocol that exactly simulates any qubit strategy in a prepare-and-measure scenario.

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 $\mathbf{y} \in \{0, \dots, N\}$ $x \in \{0, \dots, N\}$ $a \in \{0, \dots, N\}$ $b \in \{0, \dots, N\}$

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$$x \in \{0, \dots, N\}$$

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$$a \in \{0, \dots, N\}$$

$$b \in \{0, \dots, N\}$$

$$p(ab|xy) \neq \mathop{\boldsymbol{\Sigma}}_{\scriptscriptstyle{\boldsymbol{\lambda}}} \pi(\boldsymbol{\lambda}) p(a|x,\boldsymbol{\lambda}) p(b|y,\boldsymbol{\lambda})$$

$$x \in \{0, \dots, N\}$$

$$y \in \{0, \dots, N\}$$

$$m \in \{00, 01, 10, 11\}$$

$$a \in \{0, \dots, N\}$$

$$b \in \{0, \dots, N\}$$

Implications to Bell Nonlocality $x \in \{0, \dots, N\}$ $\mathbf{y} \in \{0, \dots, N\}$ $m \in \{00, 01, 10, 11\}$ $a \in \{0, \dots, N\}$ $b \in \{0, \dots, N\}$ $x \in \{0, \dots, N\}$ $\mathbf{y} \in \{0, \dots, N\}$ $a \in \{0, \dots, N\}$ $b \in \{0, \dots, N\}$

Implications to Bell Nonlocality $x \in \{0, \dots, N\}$ $y \in \{0, \dots, N\}$ BETTER $m \in \{00, 01, 10, 11\}$ $a \in \{0, \dots, N\}$ $b \in \{0, \dots, N\}$ $x \in \{0, \dots, N\}$ $\mathbf{y} \in \{0, \dots, N\}$ $\phi \rangle \in \mathcal{H}_2 \otimes \mathcal{H}_2$ $a \in \{0, \dots, N\}$ $b \in \{0, \dots, N\}$



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- 2bits + SR > 1qbit (PM scenario)
- even with general POVM measurements!

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- e.g., One bit might be enough to simulate two-qubit Bell correlations

Classical Simulation of Two-Qubit Entangled States with One Bit of Communication, arXiv:2305.19935 P. Sidajaya, A. D. Lim, B. Yu, V. Scarani

Thank you!





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